Beyond the limits of noise seismology: past and future



Three Topics

1. Dispersion measurement

2. Using the "Noise" of noise cross-correlation Function (NCF)

3. Investigating microseism PKIKP generated around Hawaii.

Ying-Nien Chen





NATURE

[October 18, 1924

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Microseisms associated with the Incidence of the South-west Monsoon.

THE late Dr. Klotz was the first to suggest a relationship between disturbed weather in the North Atlantic and the largest microseismic movements at Ottawa. The microseisms recorded by the Milne-Shaw seismograph at the Colaba Observatory during the burst of the monsoon on the west coast of the Indian Peninsula present many interesting features and indicate the possibilities of a forecast being made of the approaching monsoon at least a week ahead. The seismograph, which is installed in an underground constant temperature room, gives records remarkably free from microseisms during the cold weather period.



Lee. 1935 (particle motion analysis)

Origin: Energy Coupling between ocean wave & solid earth



Primary microseisms (PM) (Shallow water environment)



Friedrich et al., 1998

Turning noise into signal





Advantages of ambient seismic noise tomography

1.No earthquakes required! (without source uncertainties).2.Path coverage offered by station coverage.3.Quality can be improved easily by increasing data length!





An example of NCF convergence within a day (2006,101)



The mainstream of noise studies Investigating structure using ambient noise

The Flow of noise studies :



Measuring NCF dispersion without the long interstation distance limitation – a new method based on hybrid peak time matching

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Contents:

1. Motivation

2. Theoretical background

3. Method

4. Advantages of the Method

1. Motivation

High-Frequency Approximation $(r >> \lambda)$

Noise cross-correlation function (NCF)

Green's Function of the far-field surface waves

Table S1. The number of available paths for three wavelength criteria, and the selected paths used in

the inversion.

| Period | 4.0 | 6.0 | 8.0 | 12.0 | 16.0 | 20.0 |
|-----------------------------------|------|------|------|------|------|------|
| Number of paths (3λ criterion) | 3309 | 3108 | 2880 | 2370 | 1867 | 1336 |
| Qualified paths(U) | 2487 | 2428 | 2287 | 1974 | 1628 | 1300 |



Huang et al., 2015

Hybrid peak time matching method



- 1. Extending the dispersion measurement to much longer periods.
- 2. Allowing the probing of deeper structures with the same NCF dataset.

2. Theoretical background $s(\theta)$ $r^* co^{s(\theta)}$ r

Using a **"Far field "** source condition, noise source can be considered as **1**. a function of azimuth.

2. a function of the resulting delay time of a NCF.

$$NCF(\tau,\omega) \approx \sum_{\theta=0}^{n} A(\omega)^{2} \cos\left[\omega(\tau - \Delta t(\theta))\right] \qquad \Delta t(\theta) \equiv r * \cos(\theta)/c$$
$$\approx \int_{0}^{\pi} A(\omega)^{2} \cdot \cos[\omega\tau - \Phi(\theta)] d\theta \qquad \Phi(\theta) \equiv r\omega\cos(\theta)/c$$





$$NCF(\tau,\omega) \sim \Re\left[e^{-i\omega\tau} \cdot \frac{\pi}{2}A(\omega)^2 \cdot \left(J_0\left(\frac{r\omega}{c}\right) - iH_0\left(\frac{r\omega}{c}\right)\right)\right]$$

1. Peak delay time Measurement:

$$\omega \tau_{pk}(\omega) + 2N\pi = \phi \left[J_0 \left(\frac{r\omega}{c} \right) - iH_0 \left(\frac{r\omega}{c} \right) \right]$$

Phase velocity $c(T) = r/(\tau_{pk}(T))$

T:period

2. High-Frequency Approximation (r>> λ) $\omega \tau_{HF}(\omega) + 2N\pi \sim \phi \left[J_0 \left(\frac{r\omega}{c} \right) - iY_0 \left(\frac{r\omega}{c} \right) \right] = \frac{r\omega}{c} - \frac{\pi}{4}.$ Phase velocity $c(T) = r/(\tau_{HF}(T) + T/8)$

Y₀ :zero-order Bessel function of the second kind

 J_0 :zeroth order Bessel function of the first kind H_0 :zeroth order Struve function of the first kind



Comparisons of the delay time measurements predicted using different approaches. Following *Tsai* [2009], the estimated delay time is normalized by the true delay time (r/c) between stations.

3. Common measurement

If the high-frequency approximation is not valid for NCFs, the delay time measurement obtained from HF form should be

$$\omega \tau_{pk}(\omega) + 2N\pi = \phi \left[J_0\left(\frac{r\omega}{c}\right) - iH_0\left(\frac{r\omega}{c}\right) \right]$$

 $\tau_{com} = \tau_{pk}(T) + T/8$

Phase velocity $c(T) = r/(\tau_{pk}(T) + T/8)$

2. High-Frequency Approximation (r>> λ) $\omega \tau_{HF}(\omega) + 2N\pi \sim \phi \left[J_0 \left(\frac{r\omega}{c} \right) - iY_0 \left(\frac{r\omega}{c} \right) \right] = \frac{r\omega}{c} - \frac{\pi}{4}.$

Phase velocity $c(T) = r/(\tau_{HF}(T) + T/8)$







The discrepancy between the delay times given by the common and the *HF* approximation is defined as $\left(1 - \frac{\tau_{HF} + T/8}{\tau_{com}}\right)$ %. The discrepancy between them is less than 2% only when $r > 6\lambda$. Using the 2% misfit as a criterion, we argue that the *HF* approximation is only theoretically appropriate, i.e., $\tau_{HF} + \frac{T}{8} \approx \tau_{com}$, when the interstation distance between station pairs is longer than six wavelengths.

Common measurement may not be a reasonable approach in the most noise studies. However, in practice, it does provide excellent velocity measurements.





The correction of $\pi/4$ phase shift simply filters the source effects within the sensitive zone and highlight the delay time resulting from sources along station pair.

Common measurement **→** *two-station method*

3. Hybrid peak time matching method

Demonstrating how the true delay time is estimated in the proposed method by using a synthetic NCF (r = 200 km). Using the peak delay time evaluated for period 46 sec ($r/_{\lambda} = 1.08$) as an example, in the following we show the procedural to access the true delay time.





$$\omega \tau_{pk}(\omega) + 2N\pi = \phi [J_0\left(\frac{r\omega}{c}\right) - iH_0(\frac{r\omega}{c})]$$

Step 1.

Given the target period (46 sec), the x-axis is replaced with the true delay time.

$$\frac{r}{\lambda} = \frac{r}{c * T} = \frac{\tau_{true}}{T}$$

Step 2. With the true delay time (x-axis), the y-axis can be replaced with the observed delay time.

 $au_{pk} = 45.74 \ sec$ $au_{exp} = 49.37 \ sec$ $au_{true} = 49.726 \ sec$ error $\equiv |1 - \frac{\tau}{\tau_{true}}|$ error (τ_{pk}) ~8% error (τ_{exp}) ~0.6%



Period (sec)

Single frequency measurement



4. Advantages of the new method

1. The determined dispersion curves is greatly extended to include longer periods.



6 (sec) 12 (sec) 16 (sec) 20 (sec) 25 (sec) 25° (399)(259) (144)(47) 24 $r \ge 3\lambda$ 22° 120° 121° 120° 121° 122° 120° 121° 122° 122° 121° 122° 120° 121° 120° 122° 25° (624) (586) (526) (433) (504)

122°

121°

24

23

22°

121°

120°

122°

120°

This study

2. Enriching the azimuthal coverage for anisotropic tomography studies.

A comparison of path coverages at 5 selected periods obtained by applying two different approaches. In the upper panel, the black lines are the available path distributions that fit the 3-wavelength criterion; the gray lines shown in the lower panels are the paths obtained from the approach proposed in this study. The number of available paths is shown in parentheses at the top-left of each graph.

122°

120°

121°

122°

120

121°

122°

121°

120°



(a) 2D phase velocity map for 12 s Rayleigh waves derived from 399 long-distance $(r \ge 3\lambda)$ data from a traditional measurement. (b) Travel time misfits between the observations and predictions evaluated from the velocity model. The misfits of the short-distance data $(1.5\lambda \le r < 3\lambda)$ and the long-distance data are denoted by red and black dots, respectively. (c) Distribution of the misfits for the short-distance data. (d) Distribution of the misfits for the long-distance data.

Summary

- 1. We present a new method based on hybrid peak delay time matching to estimate the phase velocities at much lower frequencies of the noise cross-correlation function.
- 2. Using this method, the dispersion measurement can be extended to much longer periods, allowing the probing of deeper structures with the same noise cross-correlation function dataset.

Amplitude Asymmetry indicates the source asymmetry in ambient noise.



Friedrich at el. (1998)

2. Using the "noise" of noise cross-correlation (A New field in noise study)











Two points:

- Intrinsic noise level of NCF
 - -- (Time window) Chen et al., 2017
- Intrinsic Uncertainty of NCF
 - -- (Arbitrary lag time)

Physics Idea : Random walk (Brownian motion)

Theoretical background





Using the coherency form:

$$NCF_{j}(\omega,\tau) = \frac{1}{\overline{A^{2}(\omega)}} \sum A_{s}^{2}(\omega) \cos[\omega(\tau - \Delta t_{s})] + N(\omega) \frac{\overline{A_{\chi}^{2}(\omega)}}{\overline{A^{2}(\omega)}} \cos(\omega\tau + \phi_{j}^{avg})$$
$$\overline{NCF}^{M}(\omega,\tau) \sim \frac{\overline{As^{2}(\omega)}}{\overline{A^{2}(\omega)}} G'(\omega,\tau) + \left(\frac{INL(\omega)}{\sqrt{M}}\right) \cos(\omega\tau + \phi_{M}^{avg})$$

Intrinsic noise level



2. Using the "noise" of noise cross-correlation

Topic One : Intrinsic noise level of NCF (Chen et al., 2017)

100

0





Comparison of INL derived from model with inhomogeneous source excitation strength.



INL \rightarrow source population

Applications:

(1) Recovering Primary Microseism Signature

 $INL \rightarrow Dominant ocean wave / primary microseism Primary Microseism:$

High Population Weak Excitation



(2) Coastline Effect @ Primary microseism



NCF amplitudes \rightarrow source population & source excitation





Shallow water --> Excitation Strength (V) Complex coastline : more SPSM excitation (X)

Applications:

(2)Evaluating the exact noise level of NCF



Evaluation of INL at different lag time window of NCFs



Questions:

(1) Fluctuations of an arbitrary lag time of NCF



(2) The population of microseisms



A brief introduction to the theory

Fluctuation : Unrelated source pair

$$F_{j}(t,\omega) = \sum_{k=1}^{N} \sum_{\substack{l=1\\l\neq k}}^{N} A_{k}(\omega) A_{l}(\omega) * \cos[\omega(t - \Delta \tau_{kl}) + \Delta \phi_{klj}(\omega)]$$
(1)

The population of unrelated source pair: $N(\tau, \omega) = \sum_{k=1}^{N} \sum_{\substack{l=1 \ l \neq k}}^{N} \delta(t - \Delta \tau_{kl}).$

The mean amplitude density of unrelated source pair: $\overline{\rho}(\tau, \omega)$

$$F_j(t,\omega) = \int_{-T}^{T} N(\tau,\omega) * \bar{\rho}(\tau,\omega) * \cos[\omega(t-\tau) + \phi_j^{avg}(\tau,\omega)]d\tau$$
(2)

#############

$$F_j(t,\omega) = \int_{-T}^{T} \sqrt{N(\tau,\omega)} * \bar{\rho}(\tau,\omega) * \cos[\omega(t-\tau) + \phi_j^{avg}(\tau,\omega)] d\tau \qquad (3)$$

Source density

Source excitation

Topic two: Intrinsic Uncertainty of NCF







$$W(t,n) = 0.5 * \left[1 + \cos\left(\frac{2\pi(t-\tau)}{T^n}\right)\right] \qquad \begin{array}{l} T^n = n * period \\ \tau : \left[-T^n/2, T^n/2\right] \end{array}$$

Synthetic test (period=10 sec)



Applications: (1)Evaluating fluctuations of NCF

Predicted uncertainty: IU/\sqrt{T}

95%



Giving a quantitative description on the NCF coda quality!

Monitoring Temporal changes of seismic velocity



Brenguier et al., 2008

(2)Probing noise source properties

Intrinsic Uncertainty

On going work :

Distribution of noise source

Excitation of noise source

On the relationship between **PM** excitation and coastline geometry



Probing spatiotemporal property of solar seismic source using NCF

(Source : Convection turbulence)



Kosovichev and Alexander G. (2011)

Working with Prof. Dean-Yi Chou (NTHU) (Taiwan Oscillation Network) Active Sun Increasing small convection events

3. Investigating microseism PKIKP generated around Hawaii.

@Excitation property of body wave microseism- case study in Turkey

(b)

400

200

0

-200



0



Global Microseism Catcher Network





 $\blacktriangle IU ~ \blacktriangle II ~ \bigtriangleup G ~ \bigtriangleup GE ~ \bigstar IC$





PKIKP excitation around Hawaii







Wind stress curl (1999-2007)



deep oceans !

@ microseisms

(Pacific Decadal Oscillation: El Niño & La Niña)

@ PKIKP — Inner core structure

Thank You